Exercise 1 Charter-Penrose Diagram for Minkowski space

Consider four-dimensional Minkowski space in polar coordinates,

$$ds^2 = -dt^2 + dr^2 + r^2 d\Omega_2^2, (1)$$

$$d\Omega_2^2 = (d\theta^2 + \sin^2\theta d\phi^2). \tag{2}$$

a) Show that the coordinate change, $(r + t) = \tan(R + T)$ and $(r - t) = \tan(R - T)$ gives the metric,

$$ds^{2} = f(T, R) \left(-dT^{2} + dR^{2} + \sin^{2}(R) d\Omega_{2}^{2} \right), \tag{3}$$

and determine the function f(T, R).

- b) Give the coordinate ranges for the time coordinate T and the radial coordinate R for which Minkowski space is covered once in the new coordinates.
- c) Draw the Carter-Penrose diagram and identify i^{\pm} , i^{0} , \mathcal{I}^{\pm} .
- d) Give the paramtrization of the following geodesics through the origin R=0:
 - light rays
 - two distinct time-like geodesics
 - two distinct space-like geodesics

Draw the geodesics qualitatively in the Penrose diagram.

Exercise 2 Maxwell-Einstein equations

The aim of this exercise is to derive the solution of the Einstein-Maxwell equations of a static point-like charge. Symmetry considerations allow to consider the following ansatz for metric and electromagnetic potential,

$$ds^{2} = -f(r) dt^{2} + f(r)^{-1} dr^{2} + r^{2} d\Omega_{2}^{2},$$

$$A_{t} = \phi(r), \quad A_{r} = A_{\theta} = A_{\phi} = 0.$$
(4)

- a) Show that the current of a point-like charge Q at the origin of spherical coordinates is given by $j^{\mu}(r) = \{Q\delta(r), 0, 0, 0\}$. (Here we use the definition of the $\delta(r)$ distribution to be $\int d\theta d\phi dr \, r^2 \sin(\theta) \delta(r) f(\vec{x}) = f(\vec{0})$.)
- b) Show that the inhomogenous Maxwell equations for a generic metric tensor are given by,

$$\nabla_{\mu}F^{\mu\nu} = -4\pi j^{\nu} \quad \rightarrow \quad \frac{1}{\sqrt{g}}\partial_{\mu}(\sqrt{g}F^{\mu\nu}) = -4\pi j^{\nu} \,. \tag{5}$$

c) Solve the Maxwell equations for the ansatz (4) for r > 0 up to the integration constant. Later fix the integration constant by comparing the volume integrals of the current,

$$\int_{R^3, t=t_0} \left(\frac{1}{3!} \sqrt{g} \varepsilon_{\nu \mu_1 \mu_2 \mu_3} dx^{\mu_1} dx^{\mu_2} dx^{\mu_3} \right) j^{\nu} , \qquad (6)$$

and the left-hand side of the field equations.

Discussion 1 Reissner Nordstrøm metric

The Reissner Nordstrøm metric solves the Einstein-Maxwell equations. To show this one starts from a spherically symmetric ansatz and fixes the remaining functions using the Einstein-Maxwell equations.

The ansatz for the metric and the electromagnetic potential is given by,

$$ds^{2} = -A(r) dt^{2} + B(r) dr^{2} + r^{2} d\Omega_{2}^{2},$$
(7)

$$A_t = \phi(r) \,, \quad A_r = A_\theta = A_\phi = 0 \,.$$
 (8)

The Einstein equations are given by,

$$R_{\mu\nu} = 8\pi G_N \left(T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T_\rho^\rho \right). \tag{9}$$

The energy-momentum tensor and the non-vanishing components of the Ricci tensor are given by,

$$T_{\mu\nu} = \frac{1}{4\pi} \left(F_{\mu\rho} F_{\nu}{}^{\rho} - \frac{1}{4} g_{\mu\nu} F_{\rho\sigma} F^{\rho\sigma} \right) , \qquad (10)$$

$$R_{tt} = \frac{A''}{2B} - \frac{A'}{4B} \left(\frac{A'}{A} + \frac{B'}{B} \right) + \frac{A'}{rB}, \tag{11}$$

$$R_{rr} = -\frac{A''}{2A} + \frac{A'}{4A} \left(\frac{A'}{A} + \frac{B'}{B} \right) + \frac{B'}{rB}, \tag{12}$$

$$R_{\theta\theta} = 1 - \frac{1}{B} - \frac{r}{2B} \left(\frac{A'}{A} - \frac{B'}{B} \right) , \tag{13}$$

$$R_{\phi\phi} = \sin^2(\theta) R_{\theta\theta} \,. \tag{14}$$

The aim of this exercise will be to find the solutions of the field equations,

$$A = B^{-1} = 1 - \frac{m}{r} + \frac{q^2}{r^2}, \quad \phi = -\frac{Q}{r}, \quad q^2 = G_N Q^2.$$
 (15)

- a) Compute the field strength tensor $F_{\mu\nu}$.
- b) Compute the energy-momentum tensor of the electromagnetic field.
- c) Show that the energy-momentum tensor is trace less $T^{\mu}_{\mu}=0$.

- d) Use the explicit form of $T_{\mu\nu}$ eqns. (11) and (12) to show that $BR_{tt} + AR_{rr} = 0$ and that this implies $A(r) = B(r)^{-1} =: f(r)$.
- e) The solutions of the inhomogenous Maxwell equations is given by $\phi(r) = -Q/r$. Use this in combination with the $R_{\theta\theta}$ Einstein equations to show that f(r) is given by,

$$f(r) = 1 - \frac{m}{r} + \frac{q^2}{r^2}. (16)$$